

# Harmonic oscillator partition function from the path integral

September 11, 2018

We take the Euclidean action to be  $\dot{X}^2/2 + E^2 X^2/2$ . Directly using the energy spectrum, we can compute the partition function:

$$\begin{aligned} Z_{\text{ho}}(\beta) &= e^{-\beta E/2} + e^{-3\beta E/2} + \dots \\ &= \frac{1}{2 \sinh \beta E/2}. \end{aligned}$$

To connect to the path integral language, we write this in a more suggestive way using the infinite product formula for sinh.

$$\begin{aligned} Z_{\text{ho}}(\beta) &= \frac{1}{2 \sinh \beta E/2} = \frac{1}{\beta E} \prod_{n=1}^{\infty} \left( 1 + \frac{\beta^2 E^2}{4\pi^2 n^2} \right)^{-1} \\ &= \frac{1}{\beta E} \prod_{n=1}^{\infty} \frac{\omega_n^2}{\omega_n^2 + E^2}. \end{aligned} \tag{1}$$

In the the path integral language, the terms with  $n \neq 0$  in the infinite product come from the nonzero Matsubara modes (after dividing by the answer for the free particle with  $E = 0$ ). The only thing left to explain is why the  $n = 0$  mode gives  $1/(\beta E)$ . The path integral weight for the  $n = 0$  mode is

$$\exp -\beta \times \frac{1}{2} E^2 X_0^2.$$

Integrating this over  $X_0$  will give  $\sqrt{2\pi/\beta E^2}$ . Thus the  $1/E$  comes out correct, but the factor of  $2\pi$  and  $\beta$  needs to be fixed.

The trouble with the  $n = 0$  mode is that we cannot figure out the proper measure factor multiplying  $dX_0$  by comparing it to the free answer because the zero mode when  $E = 0$  is not Gaussian suppressed, and gives  $\infty$ . One nice way to fix this is to consider the real-time propagator

for the free particle which is given by

$$\langle x_f | e^{-iH_{\text{free}}t} | x_i \rangle = \sqrt{\frac{1}{2\pi it}} \exp -\frac{(x_f - x_i)^2}{2it}.$$

The prefactor here is fixed by the requirement that as  $t \rightarrow 0$ , the propagator must approach  $\delta(x_f - x_i)$ . To go to Euclidean signature, we have to replace  $it$  with  $\beta$ . In order to get the partition function, we have to take the trace. This involves setting  $x_f = x_i$ , which gets rid of the exponential factor. The integral over  $x$  gives the divergence proportional to the “volume of the target space”.

$$Z_{\text{free}}(\beta) = \frac{1}{\sqrt{2\pi\beta}} \times (\text{vol } X).$$

This means that

$$\begin{aligned} \left( \int dX_0 \right) \prod_{n>0} \frac{1}{\omega_n^2} &= \frac{1}{\sqrt{2\pi\beta}} \times (\text{vol } X), \quad \text{or} \\ 1 &= \frac{1}{\sqrt{2\pi\beta}} \prod_{n>0} \omega_n^2 \end{aligned}$$

We take the naive answer for the harmonic oscillator and ‘insert one’:

$$\begin{aligned} Z_{\text{ho}}(\beta) &= \sqrt{\frac{2\pi}{\beta E^2}} \prod_{n>0} \frac{1}{\omega_n^2 + E^2} \\ &= \sqrt{\frac{2\pi}{\beta E^2}} \frac{1}{\sqrt{2\pi\beta}} \prod_{n>0} \frac{\omega_n^2}{\omega_n^2 + E^2}. \end{aligned}$$

The last expression is precisely what was needed, see (1).